Whole mouse brain imaging using optical coherence tomography: reconstruction, normalization, segmentation, and comparison with diffusion MRI

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Abstract. An automated massive histology setup combined with an optical coherence tomography (OCT) microscope was used to image a total of \( n = 5 \) whole mouse brains. Each acquisition generated a dataset of thousands of OCT volumetric tiles at a sampling resolution of \( 4.9 \times 4.9 \times 6.5 \) \( \mu \)m. This paper describes techniques for reconstruction and segmentation of the sliced brains. In addition to the measured OCT optical reflectivity, a single scattering photon model was used to compute the attenuation coefficients within each tissue slice. Average mouse brain templates were generated for both the OCT reflectivity and attenuation contrasts and were used with an \( n \)-tissue segmentation algorithm. To better understand the brain tissue OCT contrast origin, one of the mouse brains was acquired using dMRI and coregistered to its corresponding assembled brain. Our results indicate that the optical reflectivity in a fiber bundle varies with its orientation, its fiber density, and the number of fiber orientations it contains. The OCT mouse brain template generation and coregistration to dMRI data demonstrate the potential of this massive histology technique to pursue cross-sectional, multimodal, and multisubject investigations of small animal brains. © 2017 Society of Photo-Optical Instrumentation Engineers (SPIE) [DOI: 10.1117/1. NPh.4.4.041501]

Keywords: massive histology; image reconstruction; optical coherence tomography; tissue segmentation; diffusion MRI; brain normalization.

Paper 17052SSPRR received Mar. 30, 2017; accepted for publication Jun. 19, 2017; published online Jul. 11, 2017.

1 Introduction

Conventional histology is widely used to explore the microstructural properties of tissue samples or to validate macroscale measurements made with other imaging modalities such as MRI. In spite of the versatility and maturity of this technique for identifying various biological tissue components, it can be labor intensive and difficult to implement when microscopic measurements over whole organs are required. For example, the BigBrain project required around 1000 h of acquisition time alone to scan a whole human brain and each tissue slice had to be manually processed to fix defects introduced by the cutting process. To address these technical limitations, an optical microscope can be combined with a motorized tissue slicing apparatus to automatize the acquisition process and thus cover large samples even with a limited field of view. Many imaging modalities can be used in such a way and are reported in the literature: scanning electron microscopy, quantitative histopathology using block-face photography, multiphoton scanning microscopy, CARS microscopy, or polarization-sensitive optical coherence tomography (PS-OCT), to name a few. In this paper, optical coherence tomography (OCT) was combined with a vibratome and a motorized stage to measure the intrinsic brain tissue contrast. This setup possesses two advantages over conventional histology: (1) no prior staining is required and (2) multiple depth measurements can be acquired simultaneously, thus allowing for \textit{in situ} slicing, thicker tissue slices, and faster image acquisition times.

In Ref. 13, it was shown that the OCT reflectivity contrast in brain tissue is mainly caused by myelinated fibers and, in a smaller proportion, to cell body densities within the tissue. This work further suggested that white matter fiber contrast depends on their orientation in relation to the direction of the microscope optical axis (MOA). The backscattering signal is higher when the fibers are orthogonal to the sampling beam rather than parallel to it. The tissue optical attenuation with depth introduces an additional spatially dependent contrast. In a review of optical properties of biological tissues, it is reported that white brain matter exhibits larger scattering coefficients and anisotropy ratio than gray brain matter. Also, A-line signal attenuation was found to be a good indicator for differentiating white and gray matter in conventional OCT. Furthermore using a PS-OCT microscope combined with a vibratome, Liu et al. showed that the signal attenuation was significantly different among the granular layer (GL), the molecular layer (ML), and the white matter fibers in \textit{in vivo} mouse cerebellum. Thus, the apparent attenuation coefficients may be used to segment the distribution of myelinated fibers within whole brains imaged with a massive histology setup using OCT as the imaging modality.

This paper describes the following. First, the tissue preparation and acquisition protocol are presented along with a description of an in-house automated high-throughput histology setup. Then, the reconstruction algorithms developed to assemble the
thousands of OCT volumetric tiles into a single volume are summarized, followed by the method used to measure the apparent optical attenuation from the OCT A-lines reflectivity signal. Four mouse brains ($n = 4$) were imaged using this technique and were coregistered to get a template of OCT reflectivity and attenuation contrasts at a downsampled isotropic resolution of 50 μm per voxel. Using these OCT brain templates, the individual brains were mapped to a common coordinate reference frame, which allowed direct comparisons of their contrasts. A tissue segmentation method that enabled the separation of myelinated axons from other cerebral tissue components was then applied on the OCT brains. The labels obtained from the OCT brain template tissue classification procedure were used as priors to segment a fifth reconstructed mouse brain. This fifth mouse brain was imaged using diffusion MRI before the automated histology acquisition, and the assembled brain was coregistered to the MRI data. Finally, the optical reflectivity and attenuation contrasts were compared with dMRI derived metrics: fractional anisotropy (FA), maximal apparent fiber density (AFD_max), number of fibers orientations within a voxel, and fiber bundle orientation. Note that this paper is an extended and revised version of a conference proceeding presented at the SPIE Photonics West-BiOS conference in February 2017.

Similar serial OCT scanners were reported previously in other publications to image rats and mice brains, as well as human brain samples. The novelty of our method lies in an improved depth-resolved signal attenuation model and the ability to create OCT brain templates using the advanced normalization tools (ANTS). This paper contributes to understanding the OCT contrast mechanism within brain tissue using diffusion MRI metrics and multimodal image registration. The presented methodology could be used to pursue multimodal investigations with multiple subjects using serial optical coherence tomography scanners.

2 Methods

2.1 Tissue Preparation

The Animal Research Ethics Committee of the Montréal Heart Institute approved all surgical procedures in accordance with the Canadian Council on Animal Care recommendations. Five C57Bl/6 mouse brains were used for this study; the brains are part of another ongoing investigation. The tissue preparation procedure followed the methodology presented by Ragan et al. Briefly, mice were anesthetized under 2% isoflurane and perfused with 20 ml phosphate buffered saline and then by a mixture of 4% paraformaldehyde with 1% gadolinium. Each mouse brain was skull extracted and embedded in 4% agarose cylindrical blocks for serial imaging.

2.2 dMRI Acquisition

One of the brains was imaged ex vivo with a standard three-dimensional (3-D) spin echo diffusion MRI sequence, using an Agilent 7 Tesla scanner equipped with 600 mT/m gradients and a custom-built 1-loop cylindrical coil. Sequence parameters were: TE = 0.021 s, TR = 0.4 s, 30 gradient-encoding directions with $b = 2079$ s/mm² and five acquisitions with $b = 0$, $\delta = 5$ ms, $\Delta = 12$ ms, gradientamplitude = 320 mT/m, FOV = 19.2 × 12.8 × 12.8 mm, and an acquisition matrix of 128 × 96 × 96 giving a resolution of 150 × 133 × 133 μm, for a total acquisition time of 36 h. The diffusion MRI data preprocessing consisted of three steps: (1) correction of the field homogeneity artefacts, (2) reduction of the Rician noise bias, and (3) volume resampling to 0.133 mm isotropic resolution.

Motion artefacts corrections were unnecessary due to the ex vivo nature of the sample. Then, in-house implementations of diffusion tensor imaging and high angular resolution diffusion imaging reconstructions were performed using the Dipy library. The FA was computed from the local diffusion tensors with a non-negative least square method. The constrained spherical deconvolution of Dipy was used to reconstruct the fiber orientation distribution functions (fODF). The principal directions of diffusion in each voxel and the AFD_max were extracted from the fODF. AFD_max is the maximal value of the fODF on the sphere, and it can be interpreted as the maximum of the apparent fiber density (AFD). Finally, the number of fiber orientations (NuFO) within a voxel was computed with the method presented by Dell’Acqua et al. and using a data-driven threshold set to 1.5 times the AFD_max values in the ventricles. The purpose of this threshold in the ventricles was to remove noisy peaks of the AFD, which are simply high-frequency peaks with low IOD amplitude. The ventricles are used because NuFO should be 0 there.

2.3 SS-OCT and Image Acquisition

An in-house automated high-throughput histology setup combined with a swept-source optical coherence tomography (SS-OCT) microscope was used to image the agaroze embedded rodent brains (Fig. 1). The SS-OCT system was reported in previous publications. The swept-source laser was operated at a central wavelength of $\lambda_0 = 1310$ nm with a tuning bandwidth of $\Delta \lambda = 100$ nm (Axsun, 1310 Swept Source Engine). A cropped Blackman apodization function was used to reduce the side lobes caused by the rectangular shape of the swept-source spectrum. The 3× microscope objective (Thorlabs, LSM04 Scan Lens) was enclosed in a 3-D printed watertight immersion chamber ending with a glass coverslip. This immersion chamber had two purposes: protect the scanning lens from the water and biological tissue debris created by the slicing process and impose a constant air/water column in the sample. The measured SS-OCT system resolutions in water were $r_{ax} = 8 \mu m$ laterally and $r_z = 10 \mu m$ axially. Its sensitivity was 98.5 dB, and its sensitivity roll-off along depth was $-0.01 dB/\mu m$.

Imaging was achieved by sequentially cutting thin tissue slices (around 200 μm) with a vibrating blade and by moving the sample under the microscope objective with a motorized stage (Zaber, T-LSR150B). At each motor position, the sampling beam was raster scanned over the objective’s field of view using galvanometric mirrors. An OCT A-line was acquired for each sampling point, thus resulting in a mosaic of volumetric OCT tiles for each tissue slice. After a slice acquisition, the sample was moved axially using a motorized jack (Thorlabs, L490MZ/M), and this process was repeated until the whole tissue was sliced and imaged. Thousands of OCT volumetric tiles were acquired, each covering a size of $\Delta x \times \Delta y \times \Delta z = 2.5 \times 2.5 \times 0.8 \text{mm}^3$ with an anisotropic sampling resolution of $4.9 \times 4.9 \times 6.5 \text{μm}^3$. The volumes were assembled using the postprocessing reconstruction method presented in Sec. 2.4.

The SS-OCT full depth range was 5 mm, but only 0.8-mm A-lines were kept for the data reconstruction. This range spanned depths between 195 to 975 μm from the zero-delay position of the OCT. This range usually comprises water and tissue. Indeed, due to variations in cutting depths, the water–tissue interface
was typically located between 395 and 570 μm from the zero-delay position. The upper and lower limits of the 0.8-mm depth range were chosen (1) to ensure that there was an overlap between consecutive tissue slices, (2) to remove the auto-correlation artefacts and internal reflections due to the coverslip, (3) to remove the areas affected by complex ambiguity artefacts, and (4) to remove deeper sample areas where the signal was dominated by multiple scattered photons or by noise.

### 2.4 Volume Reconstruction

The data reconstruction consisted of three steps: (1) tile registration, (2) XY slices stitching, and (3) Z slices stitching.

First, a displacement model of the motorized sample stage was used to estimate each tile position within the mosaic. The displacement parameters used were the tile overlaps \((O_x, O_y)\), the angle \(\phi\) between the X and Y motor axes, and the angle \(\theta\) between the laser scanning reference frame and the motorized stage reference frame [Fig. 1(b)]. Given a tile at the grid position \((i, j)\) within the mosaic, its associated motor position \(\hat{D}(i, j)\) was given by

\[
\hat{D}(i, j) = R(\theta)[p_x I + p_y R(-\phi)] [1, 0]^T,
\]

where \(I\) is the identity matrix, \(R(X)\) is the two-dimensional (2-D) rotation matrix, \((n_x, n_y)\) designates the tile shape, and

\[
p_x(i) = n_x(1 - O_x)i, \quad p_y(j) = n_y(1 - O_y)j,
\]

are the cumulative horizontal and vertical motor displacements, respectively, necessary to move the stage to the mosaic tile \((i, j)\). To estimate the displacement model parameters from the data, the translations \((dx, dy)\) between adjacent tiles within a single tissue slice were estimated using a phase-correlation-based pairwise registration method. Then, by distinguishing between the average vertical displacements \((dx_j, dy_j)\) from tiles \((i, j)\) to \((i, j + 1)\) and the average horizontal displacements \((dx_i, dy_i)\) from tiles \((i, j)\) to \((i + 1, j)\), the model parameters were estimated using

\[
\hat{\theta} = \arctan\left(\frac{dy}{dx}\right),
\]

\[
\hat{\phi} = \arctan\left(\frac{dy}{dx}\right) + \hat{\theta},
\]

\[
\hat{O}_x = \begin{cases} \frac{1}{n_x} \frac{dx}{sin(\theta)} & \text{if } \hat{\theta} = 0 \\ 1 + \frac{dy}{n_x sin(\theta)} & \text{otherwise} \end{cases}
\]

\[
\hat{O}_y = \frac{1}{n_y} \frac{dy}{sin(\phi - \theta)}.
\]

Finally, the displacement model described by the extracted parameters \((\theta, \phi, O_x, O_y)\) was used to compute the \(XY\) tile positions for each tissue slice. The advantages of using this model instead of directly computing the tile positions via image registration are that it is more robust to noise and it gives accurate position at multiple resolutions even when there is missing information. Moreover, it takes advantage of the motors’ high displacement resolution (<1 μm).

The OCT volumes were then stitched within each tissue slice using weighted average linear blending over neighboring tiles overlap areas

\[
M'(x, y) = \alpha(x, y)I(x, y) + [1 - \alpha(x, y)]M(x, y),
\]

where \(M(x, y)\) is the mosaic, \(I(x, y)\) is a tile, and \(\alpha(x, y)\) is the spatially varying blending weights (Fig. 2). These were obtained by solving the Laplace equation \(\nabla^2 \alpha(x, y) = 0\) over the overlap area with the boundary condition

\[
\delta\alpha(x, y) = \begin{cases} 1, & \text{if } (x, y) \in \Omega_i \setminus \delta\Omega_i \\ 0, & \text{otherwise} \end{cases}
\]

![Fig. 1](https://example.com/fig1.png)
where $\nabla^2$ is the Laplace operator, $\delta_\alpha(x,y)$ is the boundary of the overlap area, $\Omega_I$ stands for the domain of tile $I$, and $\delta_\Omega_I$ is its boundary. In other words, $\delta_\alpha(x, y)$ is null except when the overlap area boundary lies within the tile $I$, in which case $\delta_\alpha(x, y) = 1$. The Laplace equation was solved using the 5-stencil (in 2-D) or 7-stencil (in 3-D) finite difference approaches with a Jacobi iteration scheme. The blending weights $\alpha(x, y)$ given by the diffusion equation ensured a smooth transition between neighboring tiles. Also, the equation above can be solved in 2-D/3-D and for complex overlap geometries. Indeed, the Laplace formulation does not rely on an overlap condition; it only depends on the boundary condition. For example, when stitching tissue slices together to get a 3-D volume, the overlap geometries between adjacent slices were constrained by data masks to remove tissue slicing defects. An additional parameter $w \in [0,1]$ was used to modify the blending width within the overlap area, thus resulting in the modified blending weights $\alpha'(x,y)$,

$$
\alpha'(x,y) = \begin{cases} 
0, & \text{if } \alpha(x,y) < w/2 \\
1, & \text{if } \alpha(x,y) > 1 - w/2 \\
\frac{\alpha(x,y) - w/2}{1 - w/2}, & \text{otherwise}
\end{cases}
$$

The parameter $w$ was used to control the blending seams visibility. When $w = 0$, there is no transition between the mosaic and the new tile. The extent of the blending area increases as $w$ becomes larger. When $w = 1$, the blending is done over the whole overlap area. In other words, $w$ controls the steepness of the blending. A value of $w = 0.3$ was used for the reconstruction presented in this study. Image blending acts as a spatial compounding that reduces the speckle noise in the overlap areas only. This results in a grid-shaped artifact in the reconstructed slices as no denoising occurs inside the tiles. Using the modified blending weights $\alpha'(x,y)$, the spatial extent of the grid-shaped artifact was reduced.

The last data reconstruction step was to stitch the tissue slices together to get a complete 3-D volume. The depth translations between consecutive tissue slices were computed as follows. First, the whole OCT reflectivity tissue slices were smoothed using a 2-D $XY$ Gaussian kernel of standard deviation $k = 25 \mu m$. Then, the image curvature boundaries within each mosaic were enhanced by computing the normalized 2-D $XY$ image gradient magnitude. Using the normalized gradient maps, the cross correlation between each depth position of the $n$'th + 1 slice and the last depth position of the $n$'th slice were computed. The depth associated with the maximum image gradient cross correlation was selected as the optimal translation between the $n$'th and $n$'th + 1 slices. This procedure was repeated for all adjacent slices.

The use of the image grid for this registration process was motivated by the OCT depth-dependent signal attenuation. Indeed, the tissue contrast at the bottom of a slice is attenuated by all the structures encountered by the sampling beam. These scattering structures are removed by the slicing process, which increases the signal measured at corresponding positions in the next slice. When using the OCT reflectivity to compute the depth translation between slices, the algorithm tended to align the top tissue layers together as they exhibited similar contrasts solely due to the fact that they are not affected by signal attenuation. Instead, the normalized image gradient modulus delineates the brain and fiber tracts boundaries and is not affected by signal attenuation.

Another aspect that needs to be considered when stitching the tissue slices together is the roughness and derivatization of the water/tissue interface: the slicing process does not always result in a clear cut and the tissue deformations can introduce holes in the reconstructed data. This is taken into consideration by detecting the water/tissue interface depth for each A-line and creating a data mask to keep only the pixel under the interface in the assembled volume. Finally, the tissue slices were stitched together using the 3-D blending weights given by solving the Laplace equation $\nabla^2(\alpha(x,y,z)) = 0$ and the water/tissue masks. The 3-D Laplace equation was solved with a 7-stencil finite difference approach with a Jacobi iteration scheme. The boundary conditions were defined in the same way as for the 2-D case explained above, except that the domain of the moving volume being added was constrained by the tissue mask.

### 2.5 Optical Attenuation Estimation

Vignetting effects were compensated for in each tile prior to using the stitching algorithm. The tissue tiles containing agarose were first identified by applying the Li threshold method to separate them into background and foreground volumes. The background mosaic tiles were averaged together, and their $Z$-axis average intensity projection (AIP) was computed. This agarose AIP was used to compensate for the 2-D vignetting effect introduced by both optic misalignment and beam scanning.

As described in van Leeuwen et al., the sampling beam of a single-mode fiber reduces to a Gaussian beam, and its confocal point-spread function (PSF) adds a depth-dependent intensity contrast to each A-line. To compensate for this effect, multiple A-lines in agarose were averaged [see Fig. 3(a)]. Then, a tissue model affected by the axial component of the Gaussian beam PSF $h(z)$ was fitted on the average agarose profile to find the beam parameters: the focal plane depth $z_f$ and the effective Rayleigh length $z_R$:

$$
h(z) = \left(\frac{z - z_f}{z_R}\right)^2 + 1.0\right)^{-1}.
$$

A single PSF profile was synthesized using the extracted parameters and used to normalize each A-line in the dataset. This method assumes that the depth-dependent signal attenuation in agarose is negligible compared with the signal roll-off introduced by the microscopes axial PSF.
The signal attenuation within each A-line was then computed. The common approach to tissue attenuation estimation is to fit an exponential function on an average A-line and to report the extracted signal decay with depth obtained from this regression. The fitted equation is usually a single-scattered photon model and assumes that the attenuation coefficient is constant over the A-line scanning range. In addition to the loss of axial resolution, another drawback of this approach is the artefact that appears when the A-line traverses inhomogeneous tissues (e.g., as shown for deeply embedded fiber bundles using PS-OCT). To address these limitations, depth-resolved local tissue attenuation coefficients were estimated using the method introduced by Vermeer et al., which assumes that detected photons have experienced a single scattering event. A second assumption of this model is that the backscattering photons measured are a fixed fraction of the attenuated photons. Following this model, the attenuation at a given depth is given by

\[
\hat{\mu}(i) = \frac{1}{2\Delta} \log \left( 1 + \frac{I(i)}{\sum_{j=i+1}^{\infty} I(j)} \right),
\]

where \(\Delta\) is the axial resolution. Another assumption of this model is that most photons have been attenuated by the end of the A-line scanning range. This is not always the case, depending on the tissue being scanned. This results in artificially high attenuation values that become larger as the A-line depth increases. To relax this condition, an extended version of the Vermeer model was used [Fig. 3(b)]. The average attenuation within the tissue was computed for each A-line using the smoothed axial gradient of the signal. The average attenuation coefficient was used to extrapolate the signal below the A-line scanning range. The Vermeer model was used with these extended A-lines. The final attenuation maps were cropped to only keep values within the original scanning range. Finally, the attenuation maps were stitched together using the same mosaic positions and blending weights as for the OCT reflectivity tiles [Fig. 3(c), right].

These estimated attenuation coefficient maps were used to compute an attenuation correction bias field \(B(x, y, z)\) profile for each A-line using Beer–Lambert’s law. The original OCT tissue slices were normalized by these attenuation correction fields during the slice stitching process, which allowed for reducing signal intensity discontinuities between consecutive tissue slices [Fig. 3(c)]

\[
B(x, y, z) = \exp \left[ -2 \int_0^z \hat{\mu}(x, y, z) \, dz \right],
\]

\[
I'(x, y, z) = \frac{I(x, y, z)}{B(x, y, z)}.
\]

### 2.6 OCT Brain Template Creation

An average OCT brain template was created to be able to compare the contrasts and segmentation results of each individual brain (see Fig. 4). As will be discussed in Sec. 2.8, this OCT template was also used by the multimodal registration of the OCT brains with the dMRI data. The template creation methodology is based on existing algorithms in the literature and uses the ANTs. The normalization process was done with a subset of the assembled brains \(n = 4\). Furthermore, to get a symmetric template, the normalization also included the brain’s

![Fig. 3](https://ebooks.spiedigitallibrary.org/journals/Neurophotonics on 09 Jun 2019 Terms of Use: https://ebooks.spiedigitallibrary.org/terms-of-use)
mirror version obtained by flipping them along the midsagittal plane. Each brain was first aligned to the Allen mouse brain reference template using rigid and affine transforms and a mutual information metric. The average volume obtained by this coarse alignment was used as an initial reference template for the normalization procedure. The normalization was done iteratively, using the average template of iteration \( k \) as a reference for the \( k + 1 \) iteration. Each iteration consisted of a sequence of rigid, affine, and symmetric normalization diffeomorphic (SyN) registrations done at four scales each and using the mutual information metric. A total of \( M = 20 \) normalization iterations were done. Finally, the OCT mouse brain templates were registered to the Allen mouse brain template using a sequence of rigid, affine, and SyN registrations. For computational reasons, the template creation was done at a subsampled resolution of 50 \( \mu \text{m} \) per voxel. These OCT mouse brain templates will be publicly available on our laboratory website.

### 2.7 OCT Brain Image Segmentation

Given the observed visual contrast between white and gray matter, the Atropos algorithm, an \( n \)-tissue segmentation method developed for MRI data and part of ANTs, was used for segmenting the brain tissue from the combination of the OCT reflectivity and attenuation contrasts. The reflectivity and attenuation contrasts were combined to form a 4-D matrix, the first three dimensions being the \( X \)-, \( Y \)-, and \( Z \)-axes and the fourth dimension being a two-length vector with elements being the reflectivity and attenuation contrasts. A Markov random field with a radius of \( 1 \times 1 \times 1 \) voxel and weight \( w = 0.3 \) was used to regularize the segmented volumes. The mathematical expression of the multivariate \( n \)-tissue segmentation method is described elsewhere. All brains were registered to the Allen mouse brain common coordinate framework (CCF) prior to the tissue classification procedure, which allowed visual comparisons of the segmented tissue classes with the Allen mouse brain atlas. This segmentation procedure was applied to all individual brains and to the OCT brain template.

The segmentation algorithm was initialized using \( K \)-means clustering combined with a Gaussian mixture model (GMM). The optimal number of classes \( n \) was determined by computing the Bayesian information criterion (BIC) for varying values of \( n \). The GMM associated with the minimum BIC value was chosen. Using this criterion, the optimal number of classes was found to be \( n = 5 \). Thus, the tissue was segmented into five classes corresponding to (1) low and (2) high attenuation gray matter, white matter fibers (3) parallel and (4) orthogonal to the direction of the MOA, and (5) other white matter regions exhibiting high attenuation and reflectivity contrasts. These tissue class names were chosen based on the visual comparisons of the segmented labels with the Allen mouse brain atlas structures and based on the reflectivity and attenuation statistics per tissue label (see the results section and Fig. 7). To restrict the segmentation process to tissue voxels only, the brain and ventricles were masked prior to the Atropos algorithm application. The brain tissue was segmented with a global intensity threshold determined by the Otsu method combined with a morphological watershed algorithm. The ventricles were excluded from this segmentation process by delineating them using the ITK-SNAP tool prior to the Atropos algorithm application.

Finally, the tissue labels obtained from the OCT brain template segmentation were used by the Atropos algorithm as a priori information to guide the segmentation in other assembled brains. An advantage of this approach is a reduction in the effect of acquisition variability and artefacts on the segmentation results by imposing an a priori bias to the classification problem. This was done with the fifth mouse brain that was also acquired with dMRI, the volume was first coregistered to the symmetric OCT template. Then, the Atropos algorithm was used as described above, but by replacing the initialization step by the OCT brain template’s tissue label map and using a prior weight of \( w = 0.15 \). The average reflectivity and attenuation coefficients per tissue classes were then computed and reported in Fig. 7.

### 2.8 OCT and dMRI Coregistration and Comparison

In qualitative estimates of the attenuation coefficient maps, white matter with different orientations was apparent in images suggesting a compensation for fiber orientation (e.g., see the anterior commissure in Fig. 9). To validate this, an assembled brain was coregistered to a FA map obtained from a dMRI acquisition of the same animal (Fig. 5). Registration templates were used to align the two imaging modalities together as this method was found to be more robust than directly aligning the individual brains. To summarize the registration algorithm, the
OCT reflectivity and attenuation maps were aligned to the OCT brain templates. Then, the precomputed transforms from the OCT brain template to the Allen mouse brain were applied. Similar registration steps were done with the dMRI metrics. The FA volume was first aligned to a publicly available dMRI mouse brain template. This FA template was then coregistered to the Allen mouse brain. Thus, both the OCT brain volumes and the dMRI data were mapped to the Allen mouse CCF, which allowed voxelwise comparisons of the contrasts. All registrations were done with ANTs with a sequence of rigid, affine, and SyN transforms at four scales and using the mutual information similarity metric. The coregistration was done at a resampled isotropic resolution of 50 μm to match the resolution of the OCT brain template.

The angles between the direction of the MOA of the OCT brain and the fiber bundle orientations given by water diffusion data were computed. The OCT signal dependence on fiber orientation could then be investigated. Other MRI metrics were also compared with the reflectivity and attenuation contrasts: the FA, the AFD_max, and the number of fibers orientations within a voxel (NuFo). Plots of the average OCT reflectivity and attenuation signal as a function of these dMRI metrics were generated as follows. First, each dMRI metric was separated into 32 equidistant intervals. Then, 3-D masks indicating the spatial location in the brain of the voxels included in each dMRI metric interval were created. These masks were applied to the OCT reflectivity and attenuation maps. The plots in Fig. 8 report the average reflectivity and attenuation values for each dMRI metric value interval. The error intervals are expressed as the standard error of the mean (SEM = σ/√N), with N being the number of voxels in each interval and σ being the standard deviation of the reflectivity/attenuation values for these voxels.

3 Results

3.1 Volumes Acquisition and Reconstruction

Typically, an acquisition of a mouse brain using a 3× microscope objective generated a mosaic of shape 7 × 5 × 70 at a sampling resolution of 4.8 × 4.8 × 6.5 μm³. This represents around 320 GB of disk space. Figure 6 shows a reconstructed coronal brain tissue slice at the original microscope’s sampling resolution. Due to the large amount of data, the OCT volumes were resampled prior to reconstruction at an isotropic resolution of 25 μm per voxel, which resulted in a final brain volume size of around 300 MB. Figure 3(c) shows a horizontal section within an assembled brain at a resolution of 25 μm/voxel with three contrasts: (1) OCT reflectivity, (2) compensated OCT reflectivity, and (3) estimated signal attenuation. The compensated OCT map was obtained by dividing the OCT reflectivity signal by the cumulated attenuation for a given depth (Eq. 12). This normalization removed the signal attenuation caused by the sampling depth and by highly scattering structures traversed by the sampling beam. The remaining contrast in the compensated OCT map was due to local differences of refractive index in the brain microstructure. Indeed, the attenuation compensation increased the backscattering signal under highly scattering structures (e.g., under the corpus callosum or the arbor vitae). It also uniformized the backscattering contrast in homogeneous areas such as in the neocortex. For the white matter fibers, the backscattering contrast was higher due to the layered structure of the myelin sheaths. Also, Leahy et al. proposed that the backscattering coefficient changes based on the fiber tract orientation in relation to the direction of the incident beam. Our data corroborate this model, as illustrated by the anterior commissure olfactory and temporal limbs in Fig. 9.

The stitching algorithms were implemented in Python and were integrated as a suite of tools in a private galaxy server running in a Docker container on an 8 cores/32 GB RAM computer (64 bits). This implementation strategy was chosen to allow multiple users to launch data reconstructions simultaneously. The galaxy server can be installed on a cluster engine to provide more processing power and thus reduce computing time if needed. The reconstruction with the galaxy server took another 3 h per brain. Tests were done with the stitching algorithm on a computing cluster (called Briaüe and located at Calcul Québec) using 4 nodes of 12 cores each and 24 GB of RAM per node. The reconstruction time using the computing cluster was reduced to about 1 h for brains at the full resolution. Thus, the acquisition process using this high-throughput histology setup took between 1/2 and 1 day overall to image and reconstruct a whole mouse brain in 3-D at an isotropic resolution of 25 μm per voxel.

3.2 Attenuation Coefficients Estimation

The 3-D attenuation map resulted in a higher contrast for all white matter fibers in the brain (see Fig. 9), even those parallel to the optical axis. For example, in the anterior commissure the

Neurophotonics 041501-7 Oct-Dec 2017 • Vol. 4(4)
olfactory and temporal limbs exhibit low and high OCT reflectivity contrasts, respectively, but a high attenuation value compared with gray matter. The average OCT reflectivity ($R$) and attenuation coefficients ($\mu$) measured for each tissue class are reported in Fig. 7. To summarize, the largest attenuation coefficient was found in the fibers orthogonal to the optical axis ($24 \pm 5$ cm$^{-1}$), followed by the myelin-rich tissue ($20 \pm 3$ cm$^{-1}$), the fibers parallel to the optical axis ($20 \pm 5$ cm$^{-1}$), the high attenuation gray matter ($15 \pm 2$ cm$^{-1}$), and the low attenuation gray matter ($11 \pm 1$ cm$^{-1}$). For the reflectivity contrast, the gray matter classes and the fibers parallel to the direction of the sampling beam had similar reflectivity (around 0.23). The myelin rich tissue ($0.35 \pm 0.05$) and the fibers orthogonal to the direction of the incident beam ($0.55 \pm 0.1$) exhibited larger reflectivity values. These results were computed using the OCT template segmentation labels. Similar results were obtained with the individual brain segmentation labels. Interestingly, the histograms in Fig. 7 show that it was the combination of reflectivity and attenuation that allowed the separation of each tissue class. Thus, if the segmentation was only based on the reflectivity contrast, it would have not been possible to separate the gray matter and fibers parallel to the MOA as they all have similar reflectivity values. Similarly, a segmentation based solely on the attenuation coefficient contrast would not be able to separate the three white matter classes as they all have similar attenuation coefficients.

The largest attenuation coefficient variations were within both fiber bundle classes. One cause of this variability is that the fibers in each group are not precisely orthogonal or parallel to the direction of the MOA, thus introducing intermediary attenuation values. In addition, the Markov random field regularization and the resolution downsampling can introduce partial volume effects that mix the contrasts from different tissue components. Other sources of attenuation coefficients variations include coregistration errors, artefacts introduced by the attenuation model, tissue slicing defects, and global brain orientation variations.

### 3.3 Correlation Between dMRI and OCT

The correlations between the diffusion MRI metrics and the optical contrasts are shown in Fig. 8. These comparisons were made on a single assembled brain that was coregistered to its FA map. This map was measured for the same animal after its sacrifice and before skull extraction and histological acquisition. A threshold value of $FA = 0.4$ was used to separate the voxels into tissue and white matter tracts. This FA threshold was chosen to extract both the single-direction fiber tracts ($FA > 0.6$) and the crossing fiber areas, which are associated with FA values between 0.4 and 0.6.

The optical contrast variation with fiber bundle orientation as measured by dMRI is shown in Fig. 8(b). Consistent with observations reported previously in the literature, the OCT signal was higher for fibers orthogonal to the direction of the MOA. For example, in Fig. 6(a), the mammillothalamic tracts (red arrows) and columns of the fornix (blue arrows) are parallel to the MOA and appear darker than the surrounding tissue, as opposed to the corpus callosum and striatum (green arrows), which contain fibers parallel to the MOA. This characteristic was also observed for the attenuation coefficients. This can be explained in part by the attenuation model assumption that states that the measured backscattered signal is a fixed fraction of the attenuated light. This assumption results in a linear relationship between reflectivity and attenuation, which can be observed in the Fig. 7 scatterplot.

The multimodal comparison of OCT with diffusion MRI also revealed that reflectivity and attenuation in fibers orthogonal to the direction of the incident beam increase with FA [Fig. 8(a)]. For fibers that are parallel to the incident beam, a negative correlation is observed between these measures. This result can
be explained with partial volume effects and the orientation dependence of the OCT signal in fiber bundles. Each fiber bundle orientation population within a voxel will exhibit a different orientation-dependent OCT contrast [Fig. 8(b)]. For crossing fiber areas, low reflectivity (parallel) and high reflectivity (orthogonal) fibers are mixed and contribute to the measured OCT signal in different proportions. As FA becomes larger, the fibers within a bundle will be better aligned. If the bundle main orientation is orthogonal to the MOA, the reflectivity signal will increase as the proportion of low reflectivity parallel fibers in the bundles diminishes. This partial volume effect is consistent with the observed negative correlation of OCT contrast with the number of fibers orientations within a voxel [NuFO, Fig. 8(f)]. Some causes of FA variations in white matter are the presence of crossing fibers, fiber dispersion, fiber bundle density, or microstructural architecture of the cellular membranes. These fiber bundle characteristics should be considered when interpreting OCT contrasts in white matter.

Another source of OCT contrast variation seems to be the fiber bundle density, as shown by the negative correlation of both reflectivity and attenuation with the AFD_max [Fig. 8(e)]. The observed decrease of the OCT signal with density is counterintuitive and will require further investigation. A hypothesis to explain this effect is that the scattering anisotropy of fiber bundles increases with fiber density, in which case forward scattering becomes prominent, thus reducing the perceived reflectivity of the fiber bundle. The decrease of attenuation with density could be due to the use of an attenuation model that assumes a constant fraction of backscattered photons to attenuated photons. Future work focusing on fiber bundles exhibiting various density values is planned, including Monte Carlo simulations that can help better understand the origin of the OCT contrast.

### 3.4 Tissue Segmentation

The tissue classification was first performed on all individual brains and on the OCT brain template without using prior information to guide the segmentation. The segmentation results for the brain template are reported in Fig. 9. Based on visual inspection, the five segmented tissue classes for each brain were similar. Quantitatively, an average Dice similarity coefficient of 0.82 ± 0.03 was obtained when comparing the segmentation results of all the individual brains. This Dice score indicates that the segmentation method is accurate when similar classification initialization was used for each brain. The segmentation differences among individual samples can be caused by coregistration errors, morphological differences among the individual brains, or differences in cutting artifacts, for example.

The optical contrast dependencies on fiber orientations were used by the segmentation algorithm to distinguish between fiber tracts and myelin-rich brain tissues. Indeed, the average FA values per tissue label [Fig. 7(b)] show that the fiber bundles orthogonal (yellow) and parallel (green) to the MOA exhibit...
a larger FA signal than those in the myelin-rich tissue (red). For example, the medulla (Fig. 9, red) contains many densely packed crossing fibers that result in high reflectivity and attenuation contrasts and in a low water diffusion FA, as opposed to the spinocerebellar tract (Fig. 9, green) that was classified as a parallel fiber tract due to its remaining orientation-dependent contrast. The low FA values computed in each tissue class can be caused by misclassification errors or by crosstalk between each label. This arises from the design of the segmentation method, which only considers the combined reflectivity and attenuation contrasts distributions without prior morphological information. As a consequence, different tissue types exhibiting similar contrasts can be classified into the same class, or voxels inside a single morphological structure are classified into different labels. An example of this misclassification can be seen in Fig. 9 where part of the corpus callosum was labeled as orthogonal fibers (yellow) and as myelin-rich tissue (red).

The segmentation results reveal that another source of OCT contrast seems to be related to the density of myelinated fibers in brain tissue, as defined by the ratio of neuronal cell bodies to neurites (myelinated axons and dendrites) volume fractions. Using optical coherence microscopy, Srinivasan et al. showed that the cell bodies exhibit lower scattering coefficients than myelinated fibers, which increase the signal attenuation in myelin rich tissue. This contrast mechanism could be responsible for the distinct classification of the thalamus and hypothalamus (Fig. 9, light blue), the medulla (Fig. 9, red), and the neocortex (Fig. 9, dark blue). Indeed, the scatterplot in Fig. 7(a) reveals that each of these tissue areas is associated with increasing reflectivity and attenuation values. Similarly, using polarization-sensitive OCT, Liu et al. showed that the attenuation coefficient was significantly different among the MLs, the GLs, and the white matter regions of mice’s cerebellar slices. In histology, the granular and molecular layers are delineated based on their neuronal cell body densities, with the molecular layer being sparse and the granular layer being densely packed. This relationship between the OCT contrast and neurite density will be evaluated in future work.

4 Discussion

This paper presented a massive histology setup that combined an automated tissue slicing apparatus with an OCT microscope. An image reconstruction method adapted to the large number of volumetric tiles generated by this microscope was summarized, and the challenges were presented. The depth-resolved local tissue attenuation coefficients were estimated using a single scattering model. Average symmetric mouse brain templates were generated for both OCT reflectivity and attenuation contrasts. These templates were used to perform a multimodal registration of an assembled OCT brain with a dMRI brain acquired for the same animal. Finally, a multivariate segmentation method based on the Atropos algorithm was used to label brain tissue automatically into five components: (1) low and (2) high attenuation gray matter, (3) myelin-rich brain tissue, and white matter fibers (4) parallel and (5) orthogonal to the direction of the MOA.

The main goal of this paper was to investigate the OCT contrast to help future interpretations of serial OCT acquisitions. Thus, the dMRI data were used to better understand the OCT contrast mechanisms within brain tissue. The data showed that the OCT signal within a fiber bundle depends on its orientation,
density, and NuFO. Furthermore, the separation of gray matter into high and low attenuation groups indicates that the neurite volume fraction of brain tissue might be an additional factor influencing the OCT contrast. These hypotheses will be investigated in a future multimodal study.

The multimodal comparisons relied on the precise coregistration of the assembled OCT brain with the MRI data. Indeed, small brain structure misalignment may cause the comparison of white fiber bundles in MRI with surrounding brain tissue in the assembled OCT brain. As the AFD max, NuFO, and main diffusion orientation metrics are defined in fiber bundles only, perfect alignment of the white matter anatomical features was necessary. Missing anatomical structures (e.g., missing cerebellum) and acquisition-related deformations can introduce registration errors. To alleviate this effect, intermediary registration templates were used. The elaboration of an acquisition protocol to limit and prevent tissue deformations between each modality would be required if more subjects need to be analyzed.

The OCT brain template created for this study can be used in a number of ways. For example, it can be used in cross-sectional studies to follow the alteration of white matter in animal models of arteriosclerosis or normal aging and generate mouse brain templates at multiple aging stages. Combined with measurement of the microvasculature in the same brains using confocal microscopy or two-photon microscopy, this could give insight into the role of neurovascular impairments on the development of neuropathologies. Also, the attenuation contrast combined with local image texture analysis can be used to measure various fiber properties (e.g., orientation, density, and dispersion), which could feed tractography algorithms designed to obtain structural connectomes in ex vivo brains. Also, the reconstruction algorithm developed can stitch other types of tissue (such as aorta or heart) and other optical modalities (e.g., multiphotons microscopy, CARS microscopy, and fluorescence confocal microscopy).

The segmentation algorithm presented in this paper does not use a priori morphological information to guide the classification. It is only based on the intrinsic reflectivity and attenuation contrasts of brain tissue as measured by the OCT microscope. A limitation of this approach is that the segmentation results are highly sensitive to the initialization and to any OCT contrast variations. Another limitation of the segmentation method is that the reported tissue class names were based on visual inspection of the segmented structures when compared with the Allen mouse brain atlas. The labels were not validated using other independent measurements (e.g., with conventional histology or with an existing white matter segmentation). The segmentation method was used mainly to gain insight into the principal contrast mechanisms of OCT in brain tissue. The classification errors and variability show that naively associating OCT contrast with white matter might not be adequate, and we aimed to identify other factors that could explain the observed contrast.

Ways to improve the tissue classification method could be to (1) use additional contrasts (e.g., retardance and birefringence measured by PS-OCT or texture based measures) and (2) use an atlas-based approach to add morphological a priori information.

In the review of biological tissue optical properties by Jacques et al., the average reduced scattering coefficient $\mu'_s$...
reported for brain tissues at a wavelength of 1300 nm is 5.19 cm\(^{-1}\). This value is an average of 8 \(\mu_0\) reported in the literature,\(^ {56-58}\) with \(\mu_0\) between 2 and 8.5 cm\(^{-1}\) depending on the area probed or the measurement protocol. The attenuation coefficient values measured in the present experiment range from 11 to 23 cm\(^{-1}\). A few tissue preparation and acquisition-related aspects could explain this discrepancy. Brain perfusion, tissue extraction, agarose embedding, waiting time between tissue preparation and acquisition, or even the duration of the MRI acquisition could all have an impact on the optical properties of the tissue. For example, the brains were all perfused with gadolinium to increase their MRI signal. This could increase the tissue absorption coefficient and thus introduce a bias in the attenuation values estimated. Another indication that the MRI acquisition could have an impact on tissue optical properties was observed during the slicing process. When compared with mouse brain sliced for another ongoing study, the mouse brains that were imaged with MRI ex vivo were affected by more cutting errors (cutting plane denivelation, water–tissue interface rugosity, tissue tearing, and floating fibers). These cutting errors suggest mechanical property changes that could also impact the tissue optical properties. This tissue alteration may be caused by thermal deposition during the MRI acquisitions. Further investigation is needed to validate this hypothesis.

Some attenuation coefficient variations could also arise from the algorithm used to estimate this optical property from the A-lines data. For instance, one reason for this discrepancy could be that the measured attenuation is modulated by the agarose profile. Indeed, for PSF extraction, the agarose attenuation was considered negligible. If this is not the case, a bias is introduced in all other measured attenuation values because those are given relative to agarose attenuation. Also, the PSF was assumed to be the same for all A-line lateral positions. This hypothesis does not hold true when large fields of view are scanned or when some parts of the microscope are misaligned. In such cases, the focal plane is curved and the PSF profile intensity should vary laterally. A better PSF extraction method considering the tissue geometry, optical aberrations, and a spatially varying confocal PSF will improve the attenuation estimation results. Such a model could also use the ray transfer matrix formalism to model the microscope setup and could integrate agarose attenuation measurements for calibration. An alternative to PSF extraction is to measure the background signal in agarose or water prior to the image acquisition. Other sources of variability may come from partial volume effects, attenuation artefacts, noise sources associated with the data acquisition, and labeling errors.

The assumption that all photons have been attenuated by the end of the A-line did not always hold. Thus, the attenuation map diverged toward large values as the scan depth increased. This effect was more visible in gray matter tissue and agarose because of their lower attenuation coefficients. Also, this effect was more important when the slices exhibit large water–tissue interface depth denivelation due to cutting errors. In this work, this limitation was addressed by extrapolating the reflectivity signal under each A-line and then using this extended A-line to compute the attenuation using the Vermeer model. The extrapolation approach can introduce artefacts because it is blind to any morphological structures that are under the A-line depth scanning range. A better approach could take advantage of the underlying tissue slices revealed by the automated histology and use signal extension by interpolation instead of extrapolation. Other possible solutions are to introduce a regularization function to relax this assumption for deeper tissue\(^ {59}\) or to use a more complete tissue attenuation model that considers multiple scattering events and scattering anisotropy.\(^ {60,61}\)

Finally, some segmentation errors were caused by the attenuation artefacts introduced when the A-lines traversed a highly scattering medium (e.g., brain tissue) followed by a low scattering medium (e.g., water or agarose). This situation was encountered in ventricles and around the brain bottom side borders. This caused the attenuation to be artificially high in these transition areas, which were mistakenly considered as white matter by the segmentation algorithm. A similar artefact arose when transitioning from low to high scattering tissue (e.g., gray to white matter). These areas resulted in negative attenuation coefficients and were thus assigned a null value. These two artefacts will be taken into account by improving the tissue interaction model to consider tissue transition explicitly.

5 Conclusion

The reconstruction method developed resulted in 3-D maps of the tissue optical reflectivity and attenuation in whole rodent brains. To get an accurate estimation of the tissue attenuation coefficient, the microscope’s confocal axial PSF was estimated from the data and used to normalize each OCT A-line prior to the stitching process. An average mouse brain template was generated for both the reflectivity and attenuation contrasts, and these templates were used to segment myelinated fibers in the assembled brains. Finally, dMRI was used to get insight into the OCT contrast origin in brain tissue. The OCT signal dependence on fiber bundles microstructure and orientation, and on the neurite density in gray matter, should be considered when interpreting OCT data in brain tissue. Future work will combine attenuation and reflectivity data to measure local microstructural properties of myelinated axons, and these measures will be compared with equivalent observations from MRI acquisitions.\n
Disclosures

The authors declare no competing financial interests.

Acknowledgments

This work was supported by a Fonds de Recherche du Québec-Nature et Technologies (FRQ-NT) to J. Lefebvre and A. Castonguay, CIHR to F. Lesage, by NSERC Discovery Grant program, as well as the institutional Université de Sherbrooke Research Chair in Neuroinformatics to M. Descoteaux and by an NSERC Discovery Grant, RGPIN-2014-06089 to P. Pouliot. Computations were made on the supercomputer Briarée from the Université de Montréal, managed by Calcul Québec and Compute Canada. The operation of this supercomputer is funded by the Canada Foundation for Innovation (CFI), the ministère de l’Économie, de la science et de l’innovation du Québec (MESI), and the Fonds de recherche du Québec-Nature et technologies (FRQ-NT). The mouse brain dMRI template used in this paper was a courtesy of the Duke Center for In Vivo Microscopy. This paper is an extended and revised version of the Proc. SPIE 10070 presented at the SPIE Photonics West BiOS conference in February 2017.

References


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